

**Pauli exclusion principle:** wave function for identical fermions must be **antisymmetric** if the particle labels are exchanged

How do we tell what symmetry the isospin configurations have?  $I = 0$  or  $1$  for NN.

Use symbolic representation:  $\uparrow = \frac{1}{2}$  and  $\downarrow = -\frac{1}{2}$

The 4 configurations  $(m_1, m_2)$  are:  $(\uparrow\uparrow)$ ,  $(\uparrow\downarrow)$ ,  $(\downarrow\uparrow)$ ,  $(\downarrow\downarrow)$

$(\uparrow\uparrow)$  and  $(\downarrow\downarrow)$  are **symmetric** - exchanging the symbols (1,2) has no effect. These correspond to total isospin  $(I, m_I) = (1, 1)$  and  $(1, -1)$

$(\uparrow\downarrow)$ ,  $(\downarrow\uparrow)$  states correspond to  $m_I = 0$ , but they have **mixed symmetry**.  $\otimes$

**Solution:** make **symmetric** and **antisymmetric** combinations of the mixed states:

**symmetric:**  $(\uparrow\downarrow) + (\downarrow\uparrow) \rightarrow (\downarrow\uparrow) + (\uparrow\downarrow) \quad (I=1, m_I = 0)$

**anti - :**  $(\uparrow\downarrow) - (\downarrow\uparrow) \rightarrow (\downarrow\uparrow) - (\uparrow\downarrow) = -\{(\uparrow\downarrow) - (\downarrow\uparrow)\} \quad (I=0, m_I=0)$

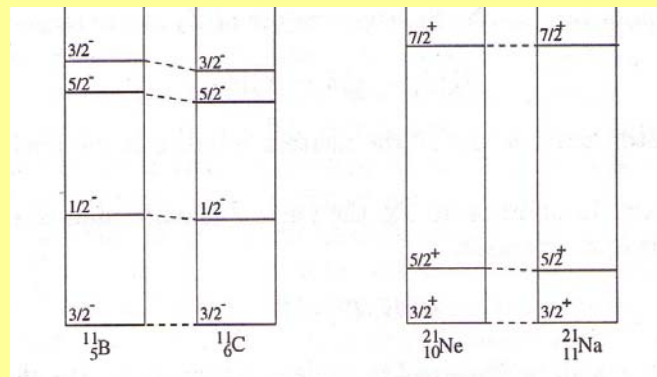
**Bottom line:**  $I = 1$  states are **symmetric**,  $I = 0$  **antisymmetric**. (Same for spin, S)  
The np system can be in a state of either  $I = 1$  or  $I = 0$  but not both, if isospin is a good quantum number.

### Isospin for nuclei:

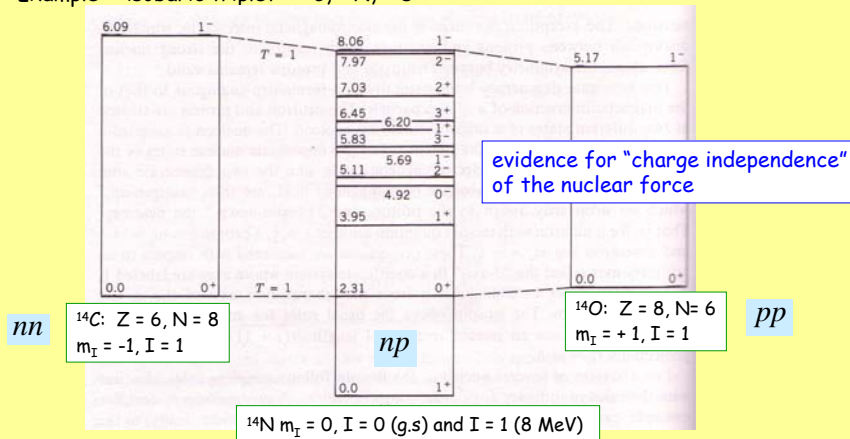
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- Nucleon:  $I = \frac{1}{2}$ ,  $m_I = \pm \frac{1}{2}$ . For a **nucleus**, by extension:  $m_I = \frac{1}{2} (Z - N)$ .
- If neutrons and protons are really "identical" as far as the strong interaction is concerned, then nuclei with the same mass number but  $(Z, N)$  interchanged ought to be very similar. These are called "**mirror nuclei**", e.g.  $^{11}\text{B}$  (5,6) and  $^{11}\text{C}$  (6,5)
- Energy spectra line up after correction for Coulomb energy difference in the ground state.  $\checkmark$

*evidence for "charge symmetry" of nuclear force*



- $pp$  and  $nn$  systems are always  $I = 1$
- $np$  system is  $(\downarrow\uparrow)$ , ie it can be partly  $I = 1$  and partly  $I = 0$
- for a **nucleus**,  $m_I = \frac{1}{2}(Z-N)$  and  $I = |m_I|$ , ie lowest energy has smallest  $I$   
(consistent with the deuteron being  $I = 0$ )
- Example: "isobaric triplet"  $^{14}\text{C}$ ,  $^{14}\text{N}$ ,  $^{14}\text{O}$ :



## Isospin selection rules for strong interactions

Consider the deuteron,  $^2\text{H} = (np)$  bound state (d)

Quantum numbers:  $m_I = 0, I = 0 \quad J^\pi = 1^+ \quad (S = 1, L = 0, \pi = (-1)^L)$

➡ How do we know it has  $I = 0$  ?

"Isospin selection rules":

The reaction: 1)  $d + d \rightarrow \gamma + ^4\text{He}$  occurs, but

isospin analysis:  $\vec{0} + \vec{0} = \vec{0} + \vec{0} \quad (I = 1 \text{ deuteron also works})$

2)  $d + d \rightarrow \pi^0 + ^4\text{He}$  does not

isospin analysis:  $\vec{0} + \vec{0} \neq \vec{1} + \vec{0} \quad (\text{only } I = 0 \text{ prevents this!})$

Bottom line:  $I$  is conserved by the strong interaction. Energy depends on  $I$  but not on  $m_I$

**Observation of the Charge Symmetry Breaking  $d + d \rightarrow {}^4\text{He} + \pi^0$  Reaction Near Threshold**E. J. Stephenson,<sup>1</sup> A. D. Bacher,<sup>1,2</sup> C. E. Allgower,<sup>1</sup> A. Gärdestig,<sup>3</sup> C. M. Lavelle,<sup>1</sup> G. A. Miller,<sup>4</sup> H. Nann,<sup>1,2</sup> J. Olmsted,<sup>1</sup> P.V. Pancella,<sup>5</sup> M. A. Pickar,<sup>6</sup> J. Rapaport,<sup>7</sup> T. Rinckel,<sup>1</sup> A. Smith,<sup>8</sup> H.M. Spinka,<sup>9</sup> and U. van Kolck<sup>10,11</sup><sup>1</sup>Indiana University Cyclotron Facility, Bloomington, Indiana 47408, USA<sup>2</sup>Department of Physics, Indiana University, Bloomington, Indiana 47405, USA<sup>3</sup>Indiana University Nuclear Theory Center, Bloomington, Indiana 47408, USA<sup>4</sup>Department of Physics, University of Washington, Seattle, Washington 98195, USA<sup>5</sup>Physics Department, Western Michigan University, Kalamazoo, Michigan 49008, USA<sup>6</sup>Department of Physics and Astronomy, Minnesota State University at Mankato, Mankato, Minnesota 56001, USA<sup>7</sup>Department of Physics and Astronomy, Ohio University, Athens, Ohio 45701, USA<sup>8</sup>Physics Department, Hillsdale College, Hillsdale, Michigan 49242, USA<sup>9</sup>Argonne National Laboratory, Argonne, Illinois 60439, USA<sup>10</sup>Department of Physics, University of Arizona, Tucson, Arizona 85721, USA<sup>11</sup>RIKEN BNL Research Center, Brookhaven National Laboratory, Upton, New York 11973, USA

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We report the first observation of the charge symmetry breaking  $d + d \rightarrow {}^4\text{He} + \pi^0$  reaction near threshold. Measurements using a magnetic channel (gated by two photons) of the  ${}^4\text{He}$  scattering angle and momentum (from time of flight) permitted reconstruction of the  $\pi^0$  "missing mass," the quantity used to separate  ${}^4\text{He} + \pi^0$  events from the continuum of double radiative capture  ${}^4\text{He} + \gamma + \gamma$  events. We measured total cross sections for neutral pion production of  $12.7 \pm 2.2$  pb at 228.5 MeV and  $15.1 \pm 3.1$  pb at 231.8 MeV. The uncertainty is dominated by statistical errors. These cross sections arise fundamentally from the down-up quark mass difference and quark electromagnetic effects that contribute in part through meson mixing (e.g.,  $\pi^0 - \eta$ ) mechanisms.

Tour-de-force experiment: <http://www.cerncourier.com/main/article/43/5/4> **$d + d \rightarrow {}^4\text{He} + \pi^0$  ???**

- isospin-forbidden reaction since  $I = 0$  for the  $d$ ,  ${}^4\text{He}$ , and  $I = 1$  for  $\pi^0$ : "textbook case"  
(technically speaking, this reaction breaks "charge symmetry" which is the symmetry under reversal of all up and down quarks in a wave function, or equivalently a quark "isospin flip". The pion wave function is  $CS - \text{odd}$ ; the others are  $CS \text{ even}$ )
- Charge symmetry is broken by the electromagnetic interaction: up-down quark mass difference, and their electric charge differences
- reaction could proceed with very low cross section compared to isospin-allowed cases, but there was never any convincing evidence published until 2003
- compare similar cross-sections at reaction threshold:

$$\begin{array}{lll} p + d \rightarrow {}^3\text{He} + \pi^0 & \sigma = 13 \mu\text{b} & (\text{Isospin allowed}) \\ d + d \rightarrow {}^4\text{He} + \pi^0 & \sigma = 13 \pm 2 \text{ pb} & (\text{forbidden, new result}) \end{array}$$

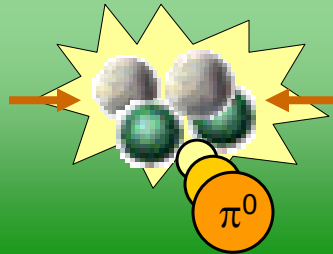
- Rough estimate of cross section ratio :

$$\sigma \sim \left( \int \psi_f V \psi_i d^3r \right)^2 \Rightarrow \frac{\sigma_{\text{allowed}}}{\sigma_{\text{forbidden}}} \sim \left( \frac{V_s}{V_{em}} \right)^2 = \left( \frac{1}{4\pi\epsilon_0 \hbar c} \right)^2 \approx 2 \times 10^4$$



Comparison of precise measurement and theory, accounting for all known CSB effects, tests our understanding of CS as a symmetry of the strong interaction

Cooler  
CSB



the search  
for  $d+d \rightarrow \alpha \pi^0$

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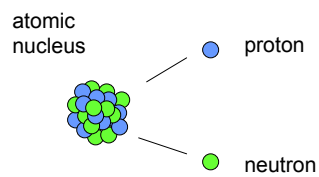
slides courtesy of Dr. E. Stephenson, Indiana University

Ed Stephenson  
Physics Colloquium  
9/24/03

full set: <http://www.iucf.indiana.edu/Experiments/COOLCSB>

## CHARGE SYMMETRY BREAKING

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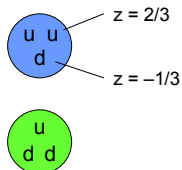


Simple notion: charge symmetry  
The proton and neutron are the same  
except for electromagnetic properties.

Isospin: the quantum number for CS  
Proton and neutron have  $I = 1/2$

But they are different:  $m_N - m_P = 1.3 \text{ MeV}$   
(The neutron decays in 887 s:  $n \rightarrow p + e^- + \bar{\nu}_e$ )

quarks  
inside  
nucleons:  
CS says  
up and  
down are  
the same  
except  
for charge

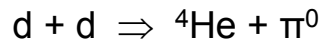


Nuclear charge symmetry breaking comes from:

- ➡ electromagnetic interactions among quarks
- ➡  $m_d > m_u$

How much does each contribute?

# Observation of the Isospin-forbidden $d+d \rightarrow {}^4\text{He} + \pi^0$ Reaction near Threshold

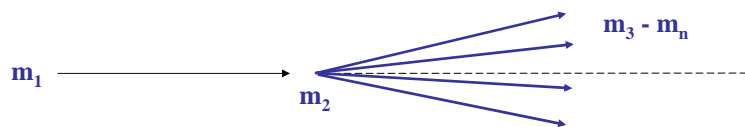


isospin: 0   0   0   1

**CHARGE SYMMETRY** says that the physics is unchanged when protons and neutrons are swapped, or when up and down quarks are swapped.

The pion wavefunction  $\psi = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d})$  is not symmetric under up-down exchange. Deuterons and helium reverse exactly. Thus, an observation of this process is also an observation of charge symmetry breaking.

## Threshold Energy:



For a reaction to occur in a fixed target experiment,  $m_1$  has to hit  $m_2$  with enough energy to make the particles in the final state. The minimum kinetic energy required is called the threshold energy:

$$T_{\text{thr}} = -Q \frac{m_1 + m_2 + \sum m_f}{2m_2}$$

$$Q = m_1 + m_2 - \sum m_f$$

Relativistic formulation! Next homework...

### Examples:

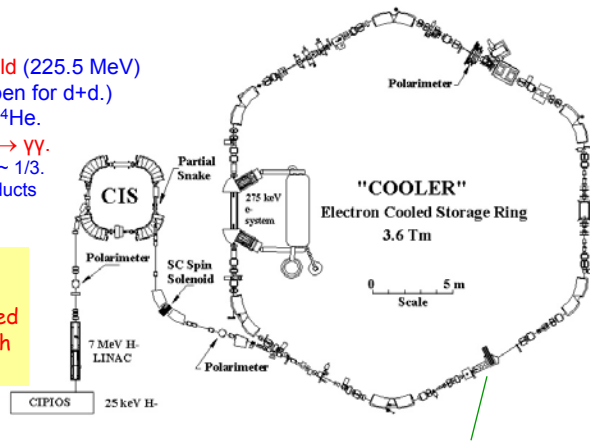


## Experimental approach:

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Search just above threshold (225.5 MeV)  
 (No other  $\pi$  channel open for d+d.)  
 Capture forward-going  $^4\text{He}$ .  
 Pb-glass arrays for  $\pi^0 \rightarrow \gamma\gamma$ .  
 Efficiency on two sides  $\sim 1/3$ .  
 Insensitive to other products  
 ( $V_{\text{beam}} = 0.51$ )

Pb-glass measures photon energy via Cerenkov light from high energy  $e^-$  produced in a 'shower' initiated by high energy photon collisions

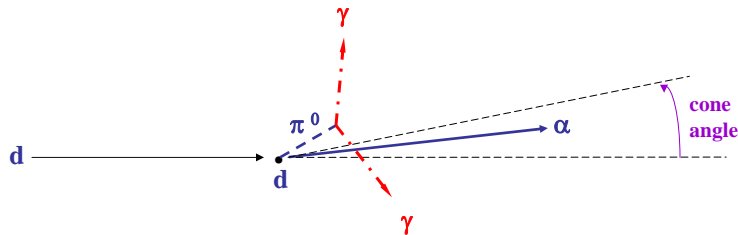


Target density =  $3.1 \times 10^{15}$   
 Stored current = 1.4 mA  
 Luminosity =  $2.7 \times 10^{31} / \text{cm}^2/\text{s}$   
 Expected rate  $\sim 5 / \text{day}$

6° bend in Cooler straight section  
 Target upstream, surrounded by Pb-glass  
 Magnetic channel to catch  $^4\text{He}$  ( $\sim 100$  MeV)  
 Reconstruct kinematics from channel time of flight and position.

## $d + d \rightarrow \alpha + \pi^0$ in the lab close to threshold:

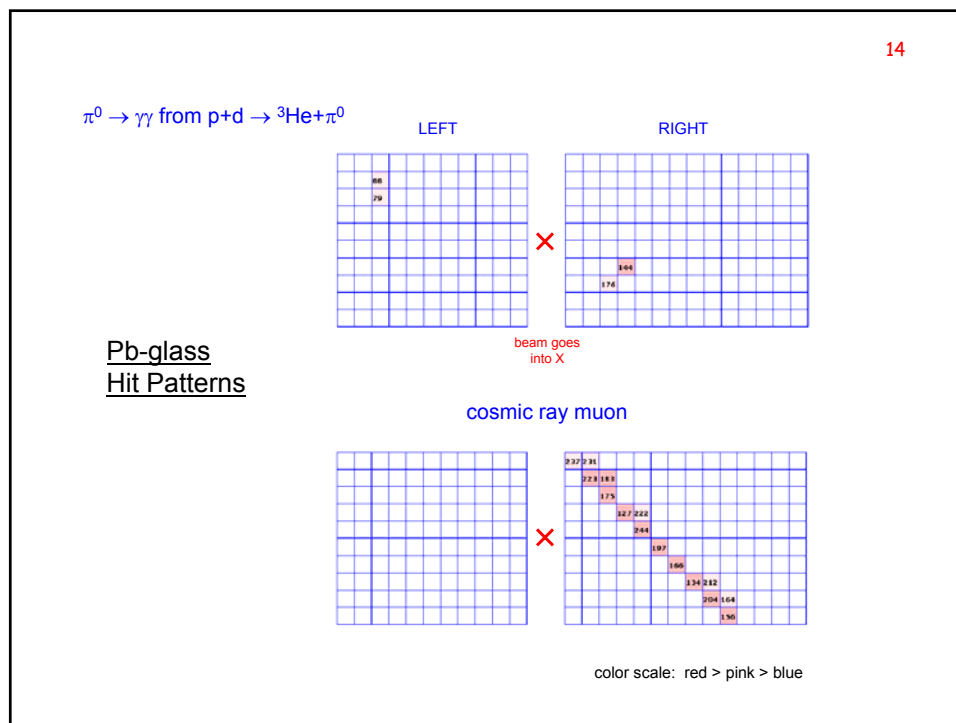
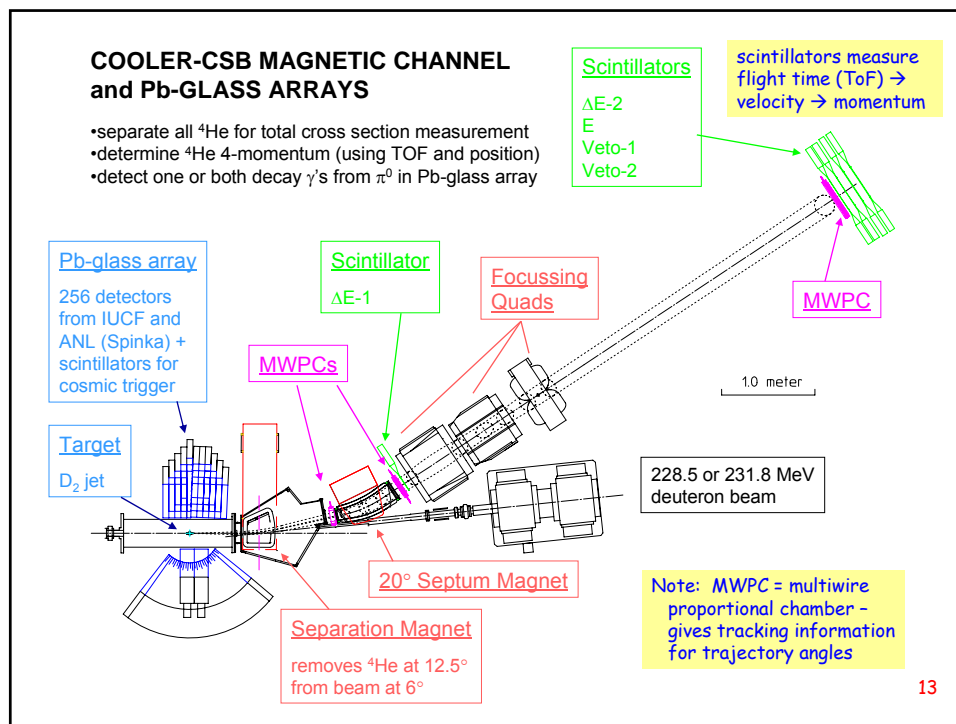
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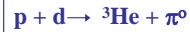
For a fixed target experiment just above threshold,

- $\alpha$  particles emerge within a narrow cone about the 0-degree line.  
 (Spectrometer with small forward acceptance will catch every  $\alpha$ .)
- low-energy  $\pi^0$  quickly decays into two photons which emerge nearly back to back in the lab.

Therefore, the apparatus must identify a forward  $\alpha$  in coincidence with two photons that have a large opening angle between them.



`Missing Mass' measured with proton beam:

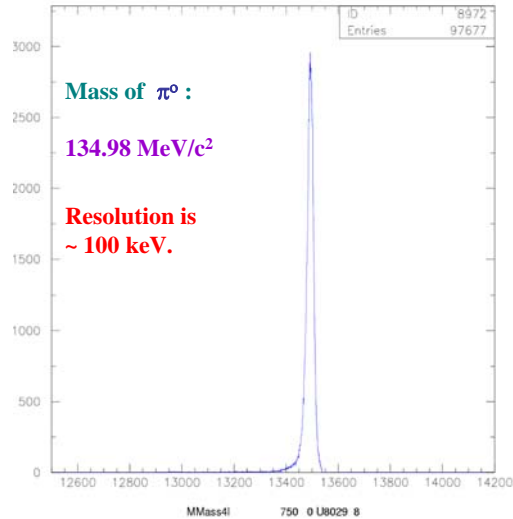


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conservation of energy:

$$W \equiv E_p + E_d - E({}^3\text{He}) = m_\pi$$

- $E_p$  from beam energy
- deuteron at rest in target
- $E({}^3\text{He})$  from energy and momentum measured with the magnetic channel
- calculate  $W$  from data, should find a peak at the pion mass for reaction at threshold.
- then check in Pb glass array to see if pion was observed

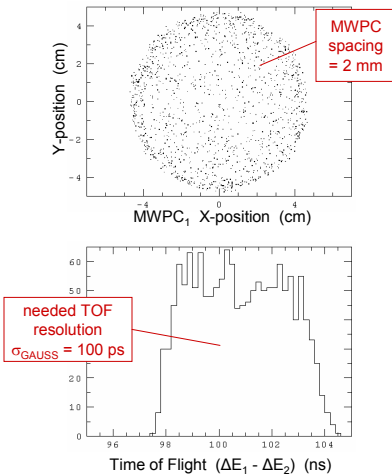


Mass ( $\text{MeV}/c^2 * 100$ )

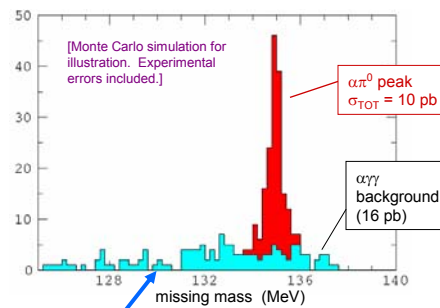
## SEPARATION OF $\alpha\pi^0$ AND $\alpha\gamma\gamma$ EVENTS

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IDEA: Calculate missing mass from the four-momentum measured in the magnetic channel, using time-of-flight for z-axis momentum and MWPC X and Y for transverse momentum. Should see a peak for  $\alpha\pi^0$  reaction and a broad background from  $\alpha\gamma\gamma$



Need very good resolution so that the peak is detectable!

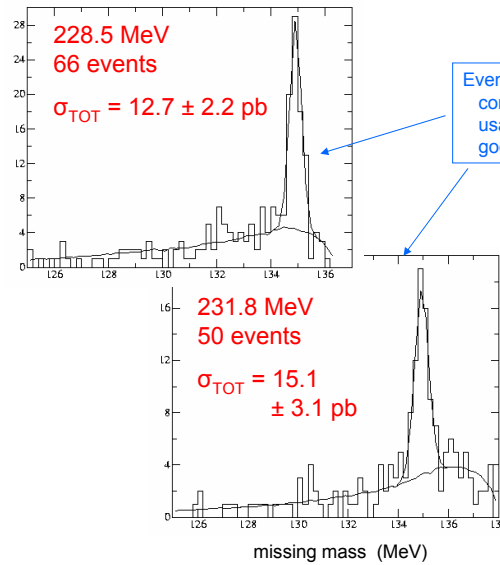


Background:  $d + d \rightarrow \alpha + 2\gamma$



**RESULTS** (measured at two different beam energies)

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Events in these spectra must satisfy:  
correct pulse height in channel scintillators  
usable wire chamber signals  
good Pb-glass pulse height and timing

First ever convincing  
observation of both the  
 $\alpha\pi^0$  and  $\alpha\gamma\gamma$  reactions!

Peaks give the correct  
 $\pi^0$  mass with 60 keV  
error. ✓

Bottom line: time to revise all the textbooks!!!